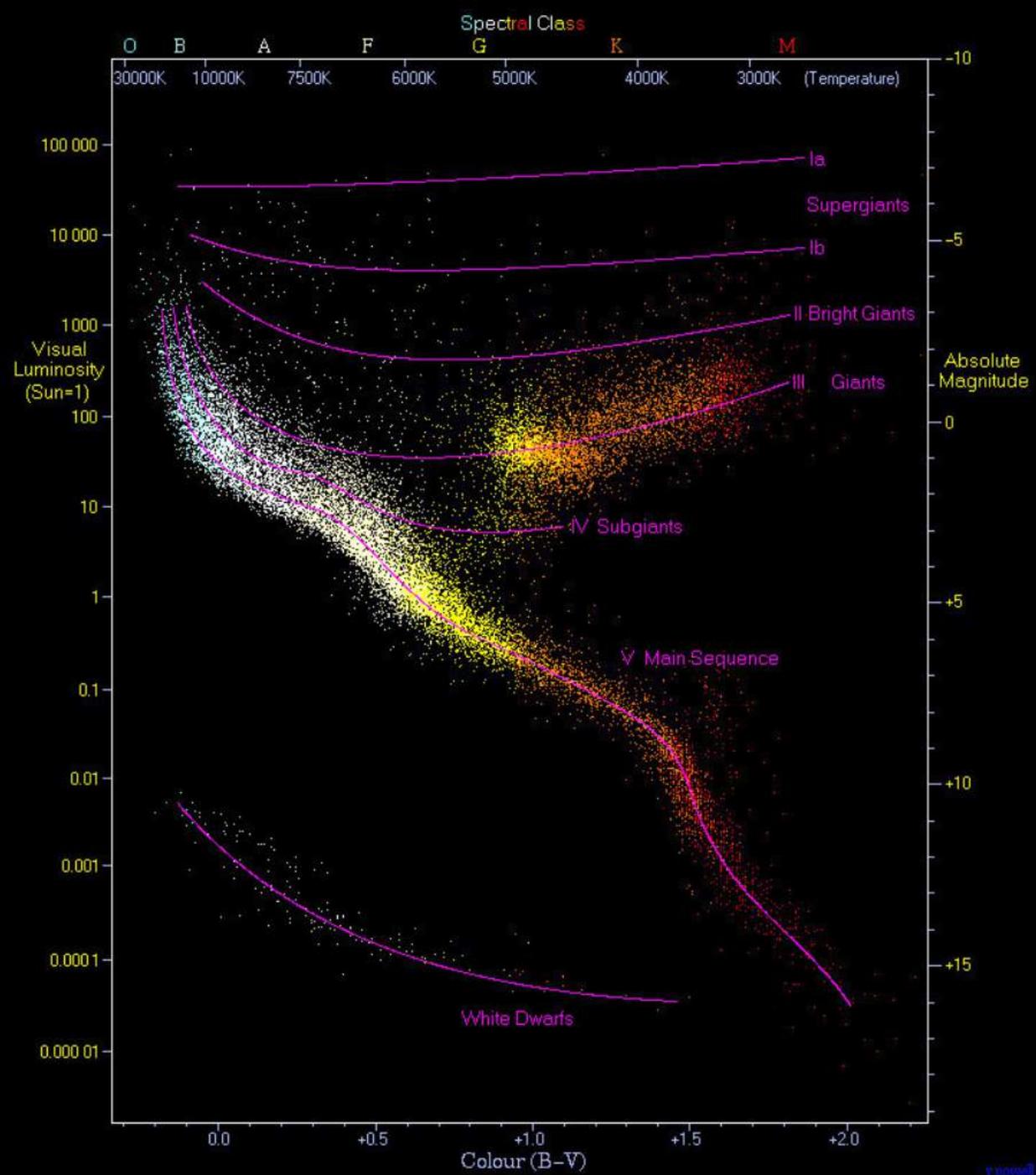


Stellar Evolution
of low and massive stars



How long do stars stay on the Main Sequence?

- Stars spend most of their lives on the Main Sequence and fuse H to He
- Lifetime depends on birth mass:

$$\tau_{MS} = \frac{E_{nuclear}}{L} \propto \frac{M}{M^3} = M^{-2}$$

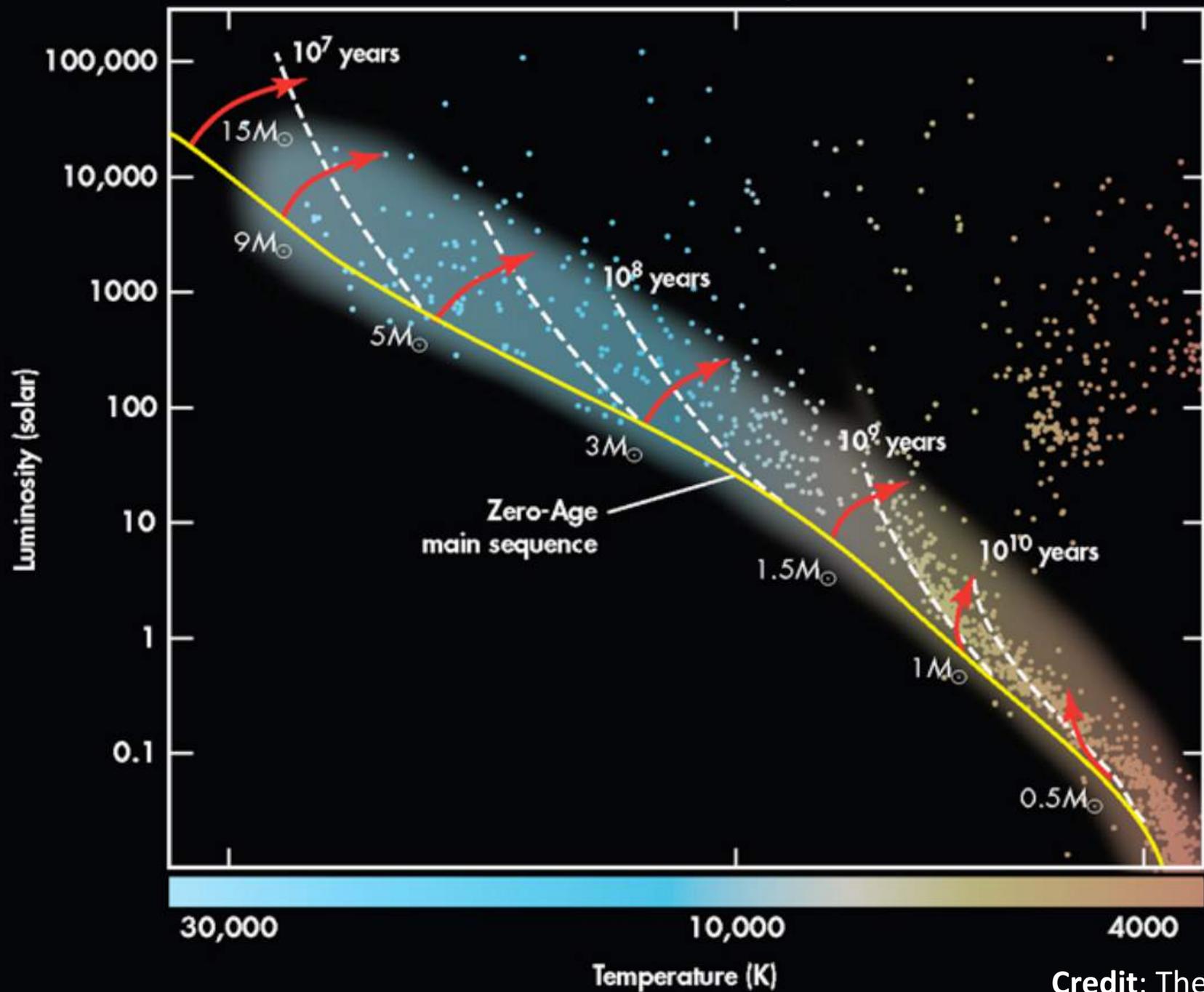
Mass-Luminosity-Relation:

$$L \sim M^3$$

→ larger mass results in shorter lifetime on MS

Mass/ M_{\odot}	MS lifetime (yrs)
0,10	$2 \cdot 10^{12}$
0,50	$2 \cdot 10^{11}$
0,75	$3 \cdot 10^{10}$
1,0	$1 \cdot 10^{10}$
1,5	$2 \cdot 10^9$
3	$2 \cdot 10^8$
5	$7 \cdot 10^7$
10	$2 \cdot 10^7$
15	$1 \cdot 10^7$
25	$7 \cdot 10^6$
60	$3,4 \cdot 10^6$

Lifetime on Main Sequence as a function of mass

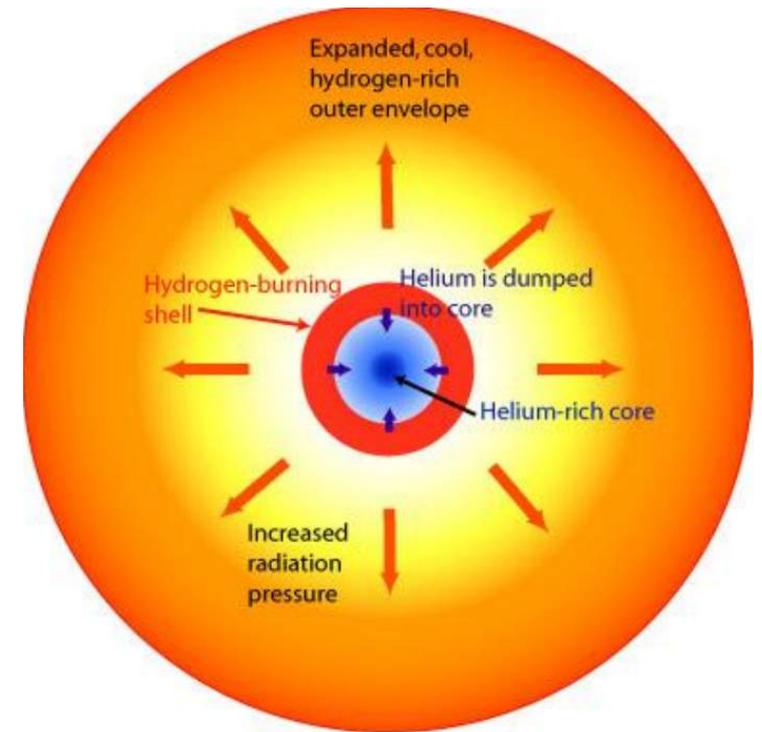


Credit: The McGraw-Hill Company

Low Mass Evolution for $M > 8 M_{\odot}$

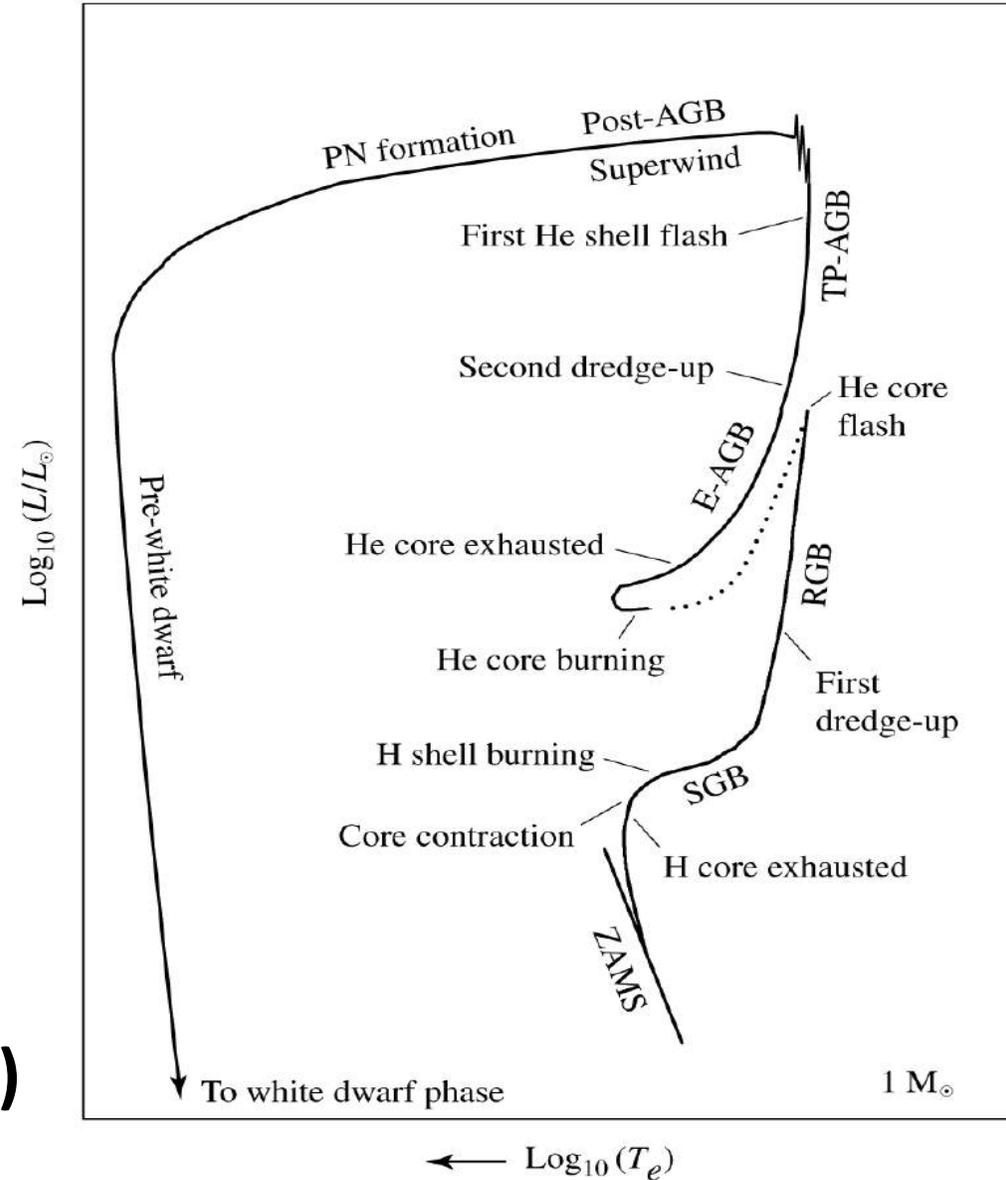
- Hydrogen in the core is exhausted
→ energy production via pp cycle stops
- T high enough for H shell burning
→ Helium core grows
→ HRD: star moves to right (cooler T)
- When reaching the Chandrasekhar limit the He-core contracts
→ the envelope expands and the star cools down
→ **Sub Giant Branch (SGB)**

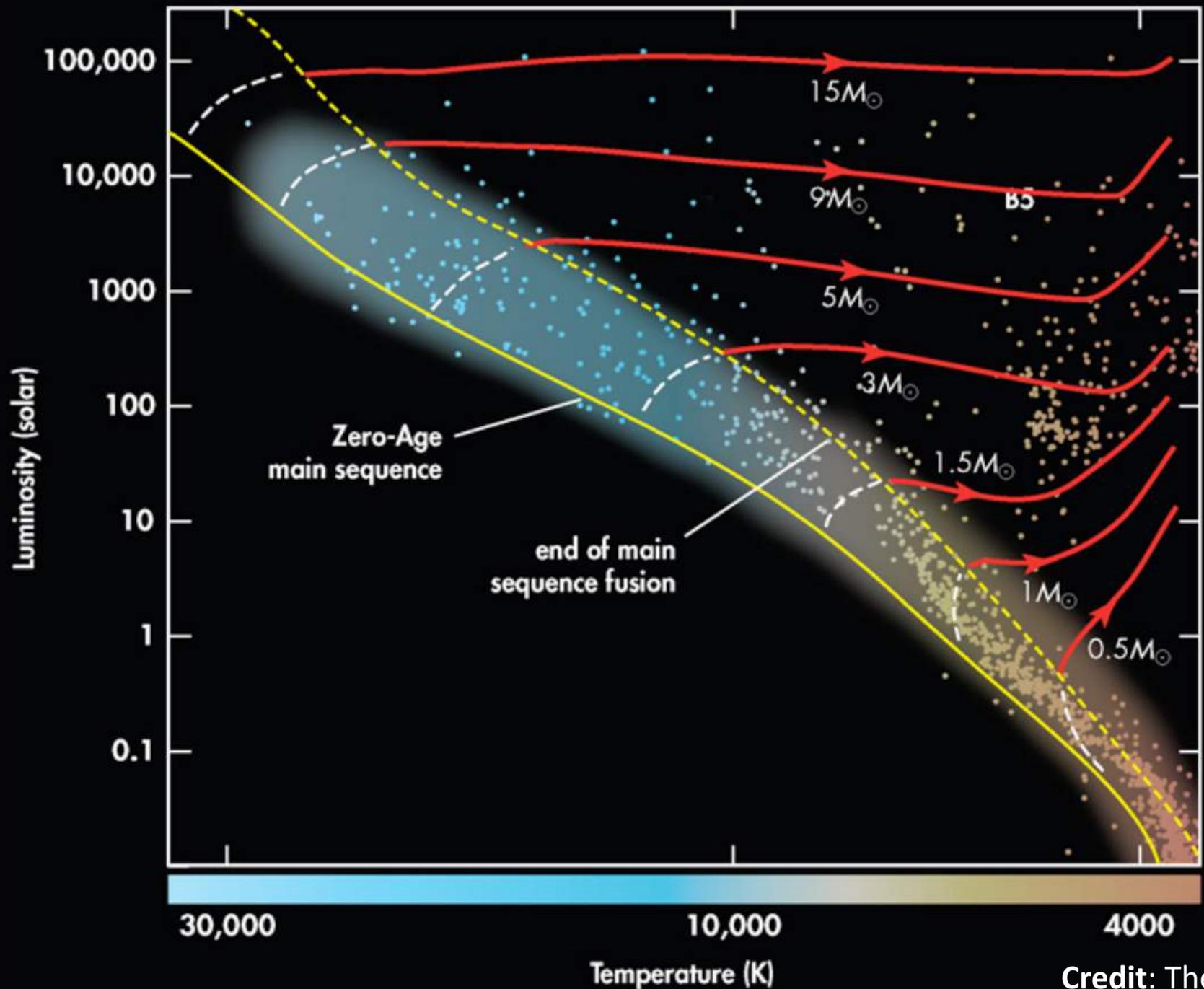
Chandrasekhar-limit: $1,4 M_{\odot}$



Low Mass Evolution for $M > 8 M_{\odot}$

- $T \sim 5000 K$: opacity of the envelope increase
→ Convection sets in
→ Luminosity increases dramatically
- Core continues to collapse
→ L and P in shell increases
→ outer layers become convective
- Star is not longer in hydrostatic equilibrium
→ envelope expands and cools down
- HRD: star moves up the **Red Giant Branch (RGB)**

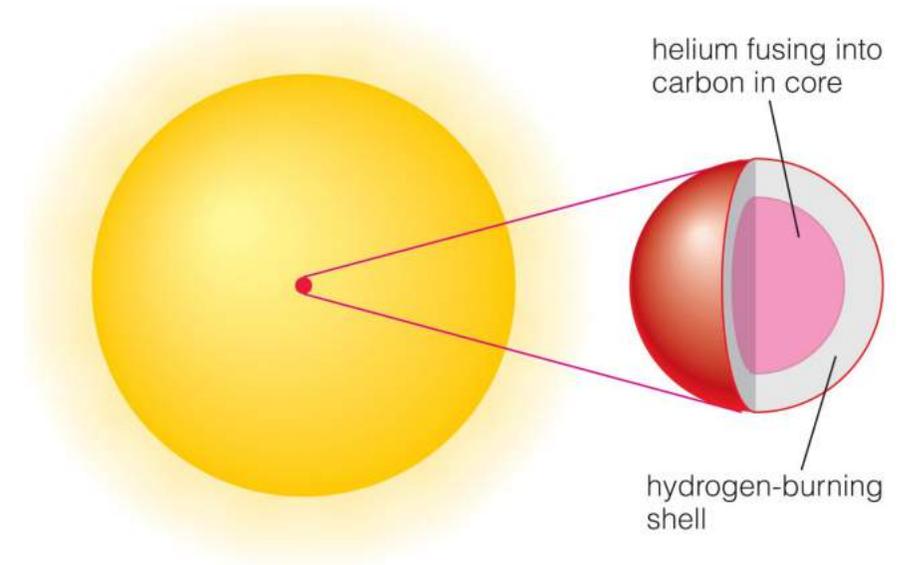
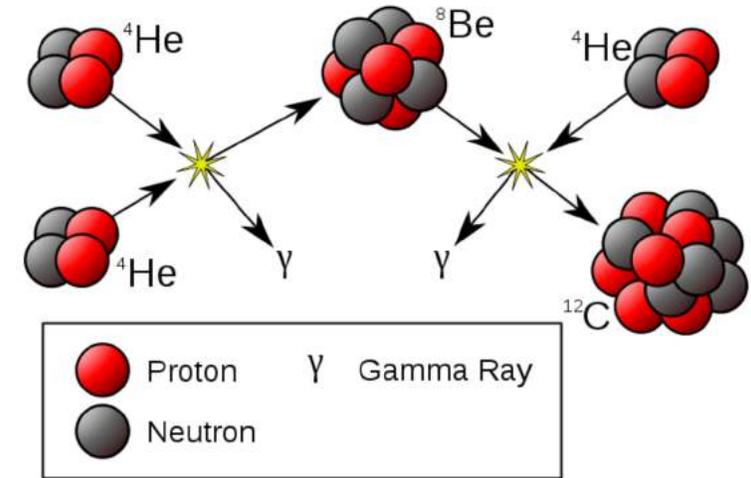


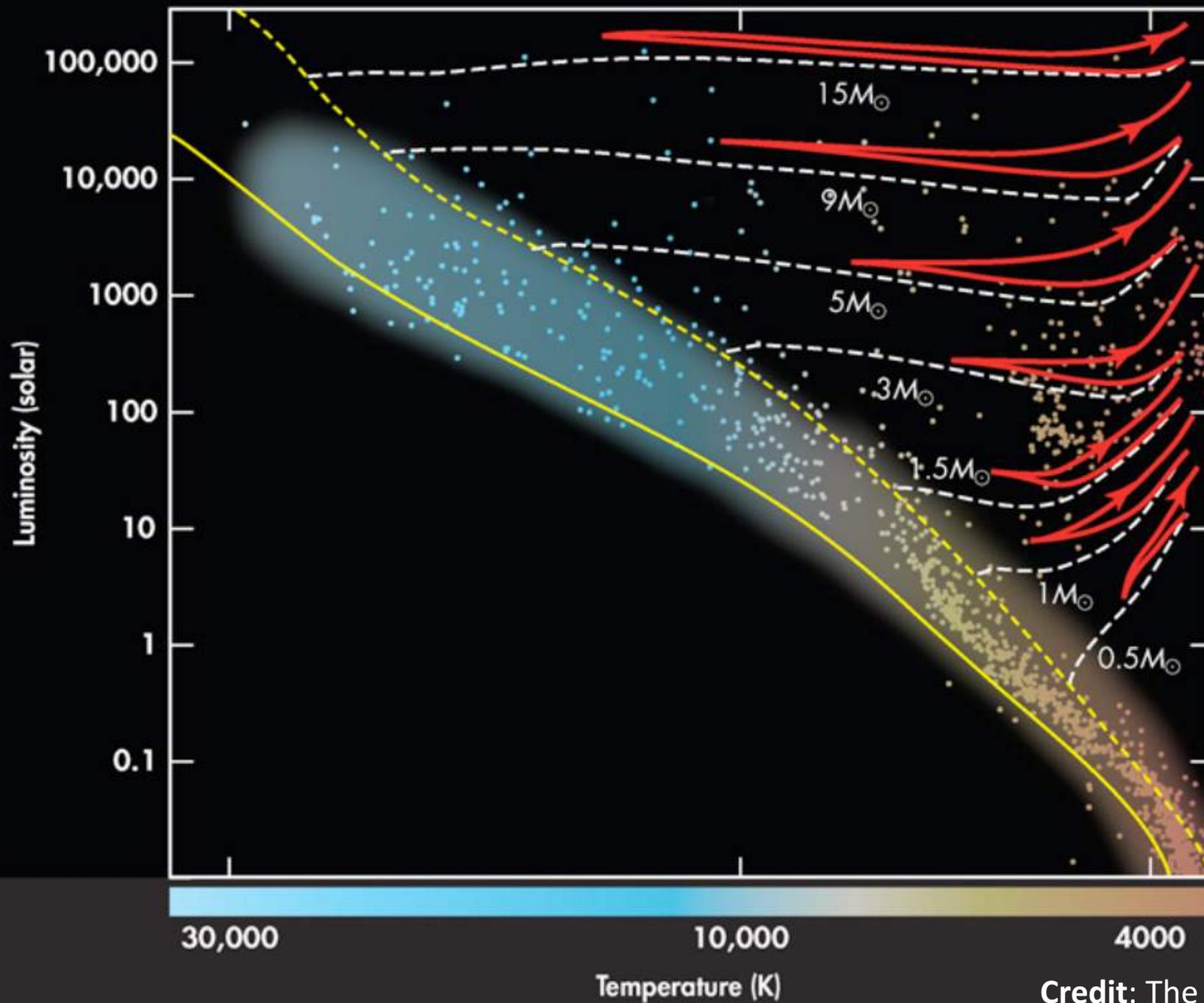


Credit: The McGraw-Hill Company

When does Helium ignition start?

- End of RGB: T_{core} is high enough for triple- α process
→ core must be degenerate for high densities
- Nuclear runaway: **Helium flash**
→ energy removes electron degeneracy
- He core burning and H shell burning
- Star starts to fuse He to C in its core
→ Lifetime as red giant $\sim 10\%$ of MS lifetime

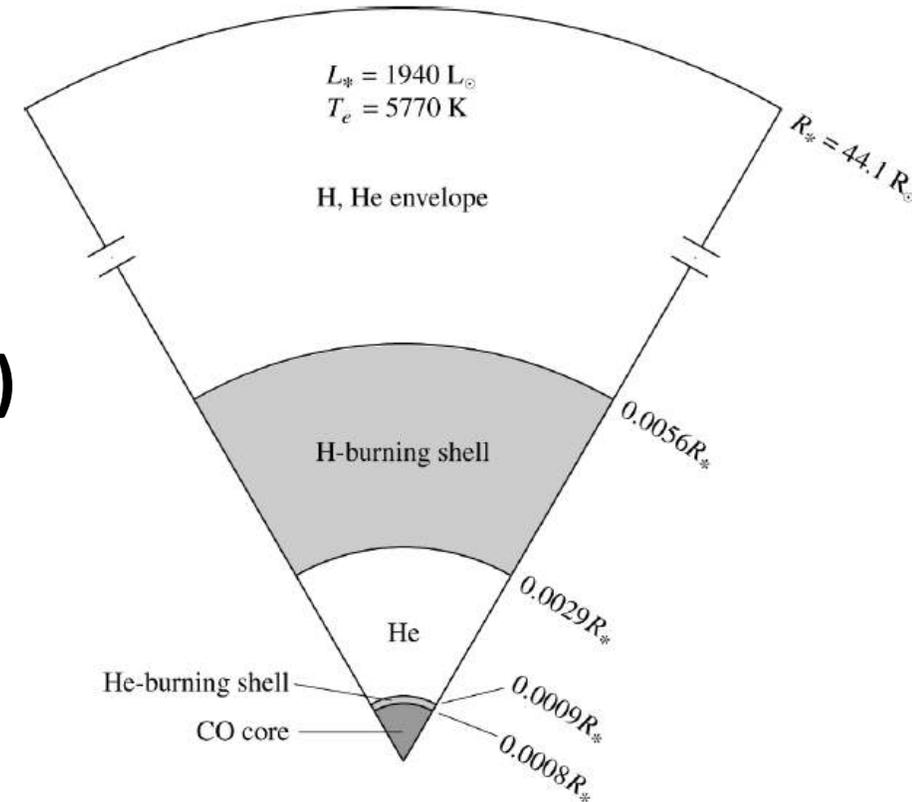




Credit: The McGraw-Hill Company

Late Burning Phase for $M < 8 M_{\odot}$

- Core runs out of He, it stops producing energy
→ begins to collapse again
- Inert C core, He and H shell burning
- Star moves up the **Asymptotic Giant Branch (AGB)**
- Shell burning occurs not simultaneously
→ changing thermal pulses
→ outer layers of the star are ejected
- Mass loss increases, until the entire envelope is ejected

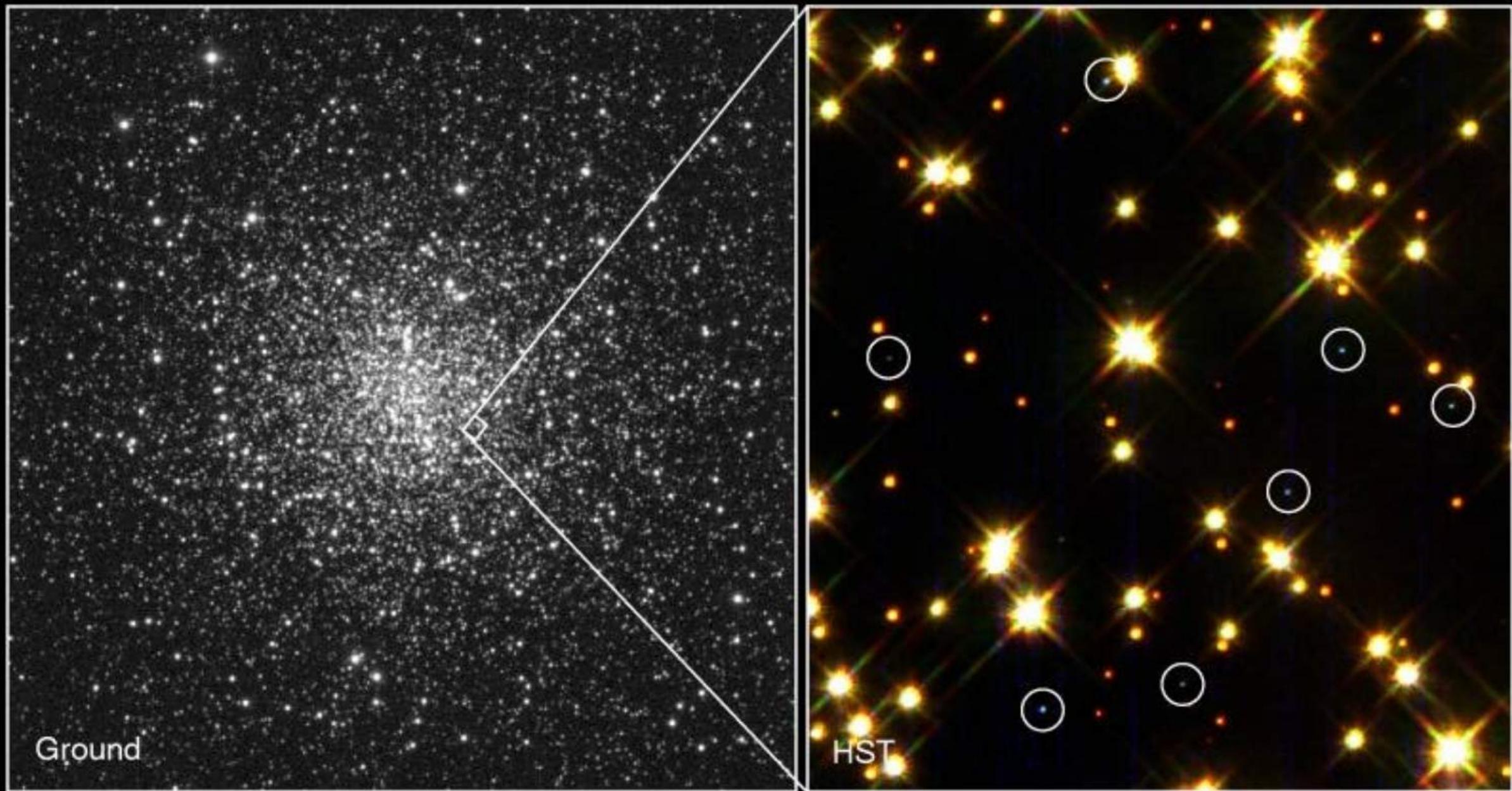


Final Stage for Low Mass Stars with $M < 8 M_{\odot}$

- Star is not massive enough to ignite C core
- Ejected shells with hot exposed cores
→ **planetary nebula phase**
- C/O white dwarf remains
- Some stars at the higher end of this phase can burn C
→ O Ne Mg white dwarfs



Phase 1 M_{\odot}	τ (yrs)
Main Sequence	$9 \cdot 10^9$
Subgiant	$3 \cdot 10^9$
Red giant	$1 \cdot 10^9$
AGB evolution	$\sim 5 \cdot 10^6$
PN	$\sim 1 \cdot 10^5$
WD cooling	$> 8 \cdot 10^9$



Ground

HST

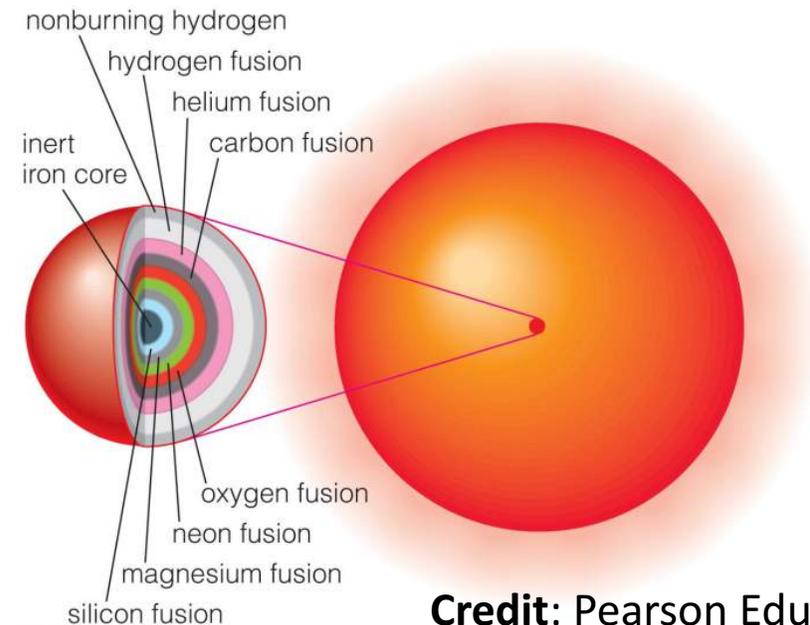
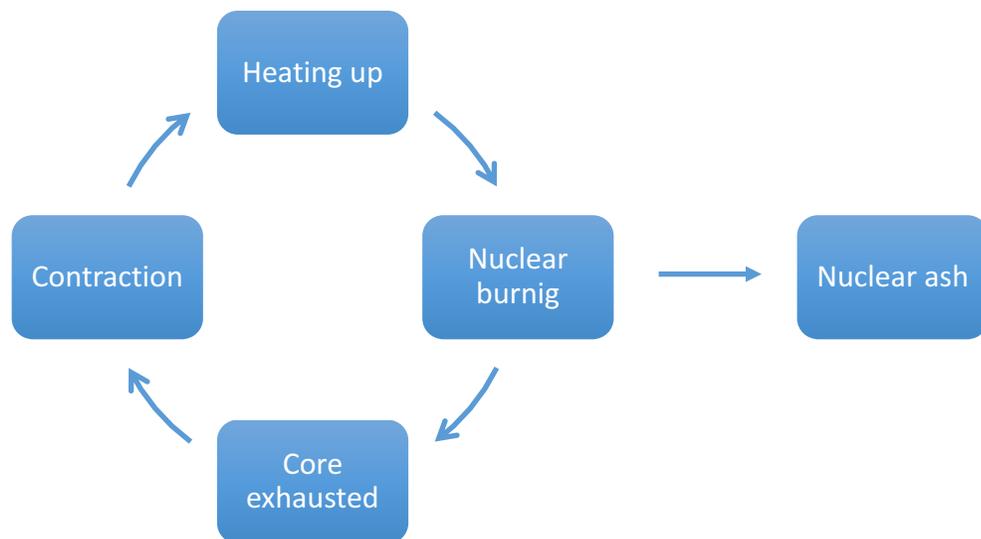
White Dwarf Stars in M4

HST · WFPC2

PRC95-32 · ST ScI OPO · August 28, 1995 · H. Bond (ST ScI), NASA

High Mass Evolution for $M > \sim 8 M_{\odot}$

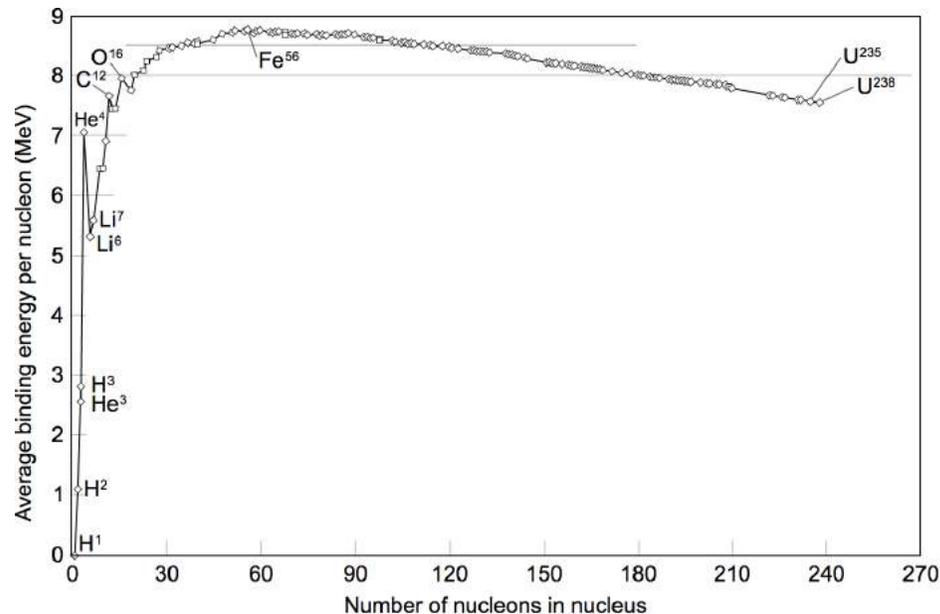
- Collapsing He core reaches T for triple- α process
→ He burning begins when core is not-degenerate (lower density)
→ no helium flash
- He burning stops when core is converted to C and O
→ core begins to collapse and He burning shell ignites outside the core
- This pattern of core ignition and shell ignition continues
→ lead to a „onion-skin“ layering of nuclei



Credit: Pearson Education Inc.

Later Burning Phases

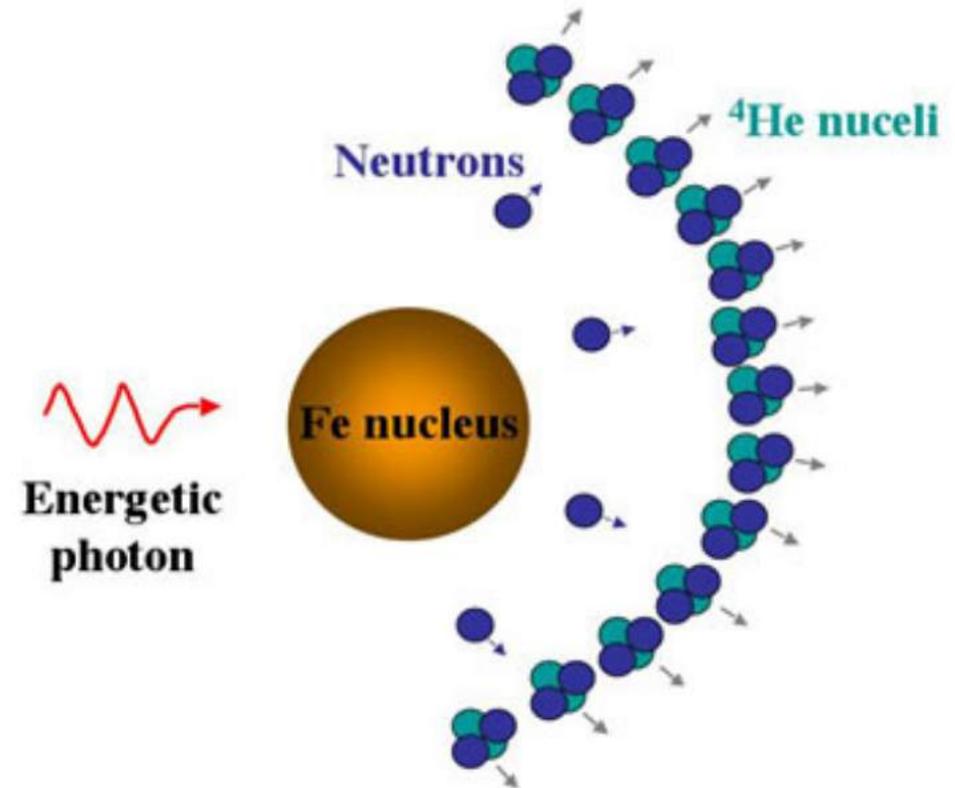
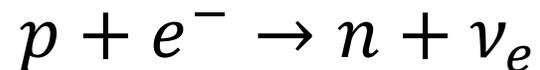
- Heavier elements are built up until Fe core is formed
 → The core is surrounded by a series of shells at lower T and lower ρ
- Nucleosynthesis of elements above He is less efficient
 → the reactions occur at greater rates so radiation pressure can balance gravity



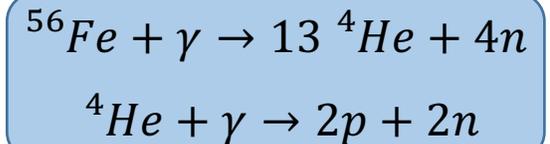
Phase	Required Temperature [K]	$15 M_{\odot}$ τ [yrs]	$25 M_{\odot}$ τ [yrs]
Hydrogen (MS)	$4 \cdot 10^7$	$10 \cdot 10^6$	$7 \cdot 10^6$
Helium	$2 \cdot 10^8$	$2 \cdot 10^6$	$5 \cdot 10^5$
Carbon	$6 \cdot 10^8$	2000	600
Neon	$1,2 \cdot 10^9$	0,7	1
Oxygen	$1,5 \cdot 10^9$	2,6	0,5
Silicon	$2,7 \cdot 10^9$	$5 \cdot 10^{-2}$	$3 \cdot 10^{-3}$

Core Collapse!

- Fe-core reaches $1,4 M_{\odot}$ (Chandrasekhar-limit)
→ degeneracy pressure can no longer resist its own gravity
- Collapse begins
→ Temperature rises, but Fe can't fuse
- Fe nuclei breaks down because of **photodesintegration**
- Inverse β -decay occurs, requires 1,3 MeV for each reaction

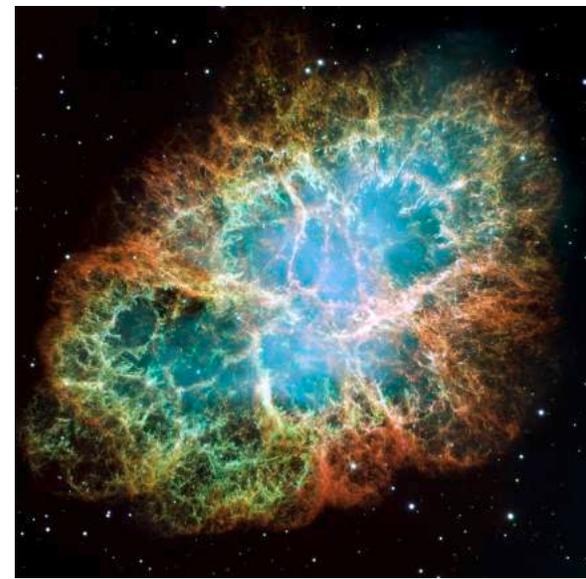


Credit: Swinburne University of Technology



Supernova

- Core collapses from under the rest of the star
→ outer layers start to fall inwards and hit surface of core
- Outer layers „bounce“ off and are again ejected out
→ creates shock wave which pushes back the envelope
- shock wave smashes outward through the star – $v \sim 0,1c$
→ explosive nuclear fusion takes place behind the shock (r-process)
→ shock heats and expands the layers of the star
- The light increases as the surface area of the ball of gas increases



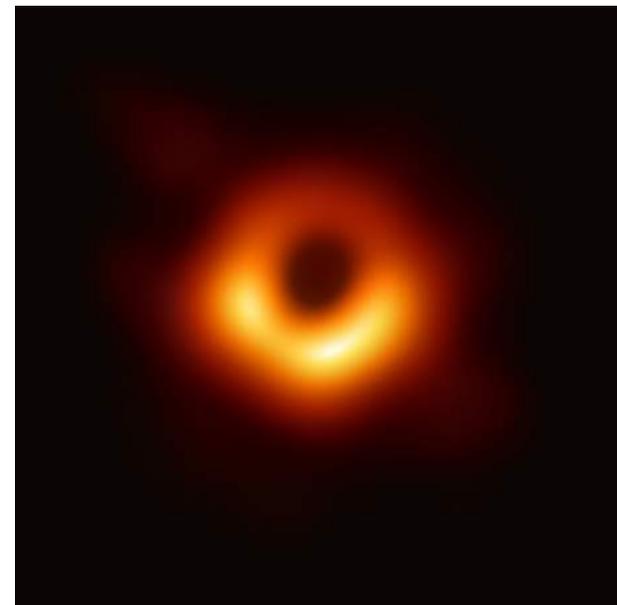
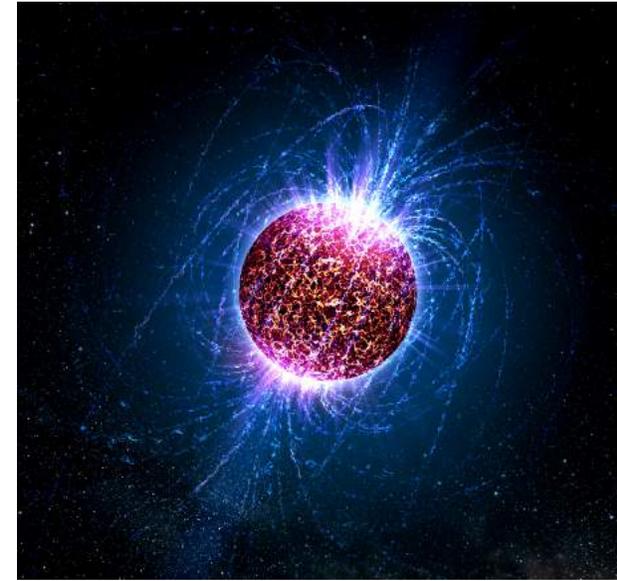
Neutron Star or Black Hole?

NS with $M < 3 M_{\odot}$:

- Neutron degeneracy sets in → resists gravitational pressure
- Whole core is like one giant atomic nucleus
→ Neutron star has been born

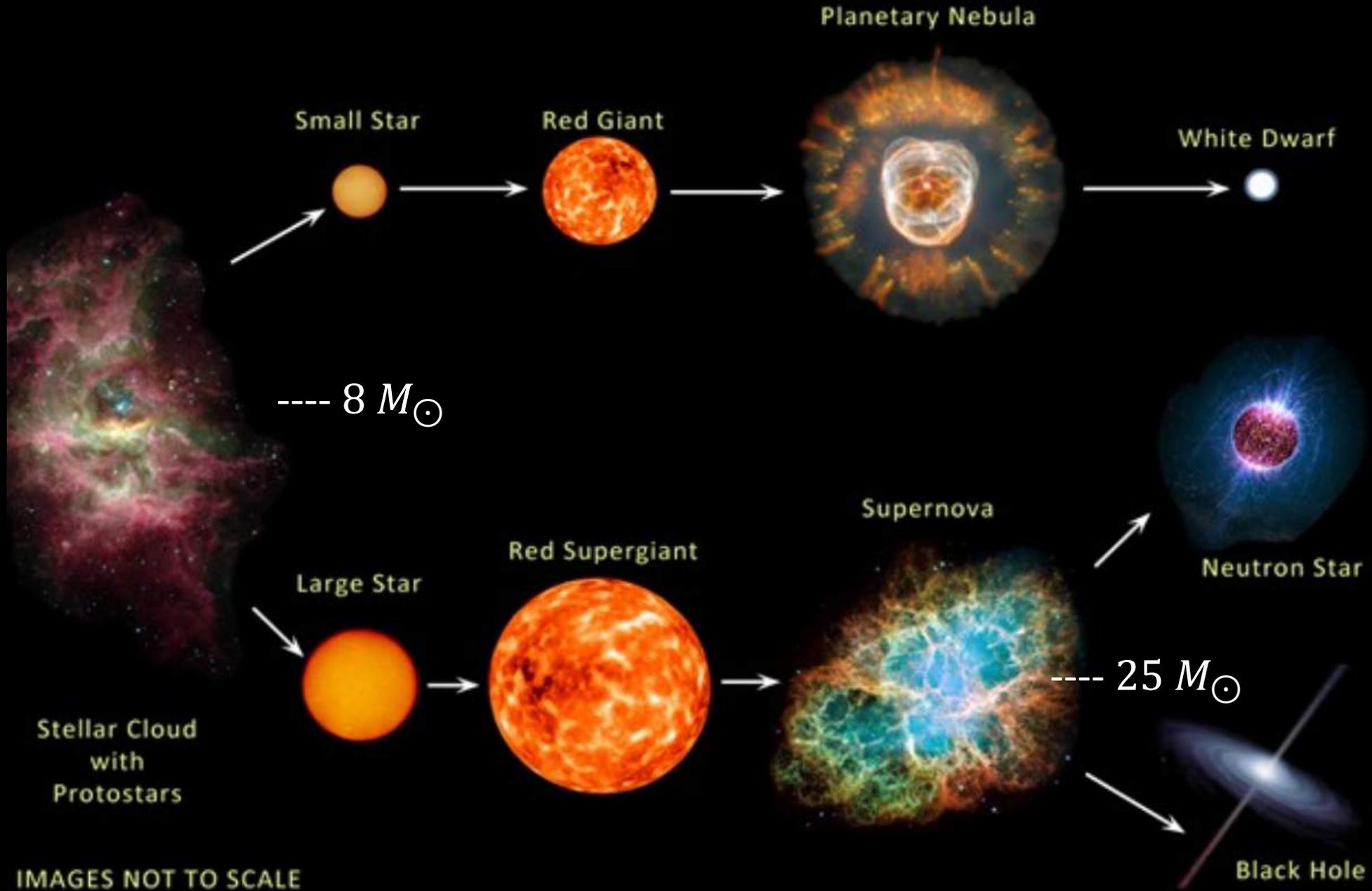
NS with $M > 3 M_{\odot}$:

- Mass of NS reaches Tolman–Oppenheimer–Volkoff limit
→ corresponding to the Chandrasekhar mass
- neutron degeneracy pressure is unable to balance self-gravity
- Upper limit is less than $\sim 3 M_{\odot}$
→ above this limit, NS will collapse to a black hole

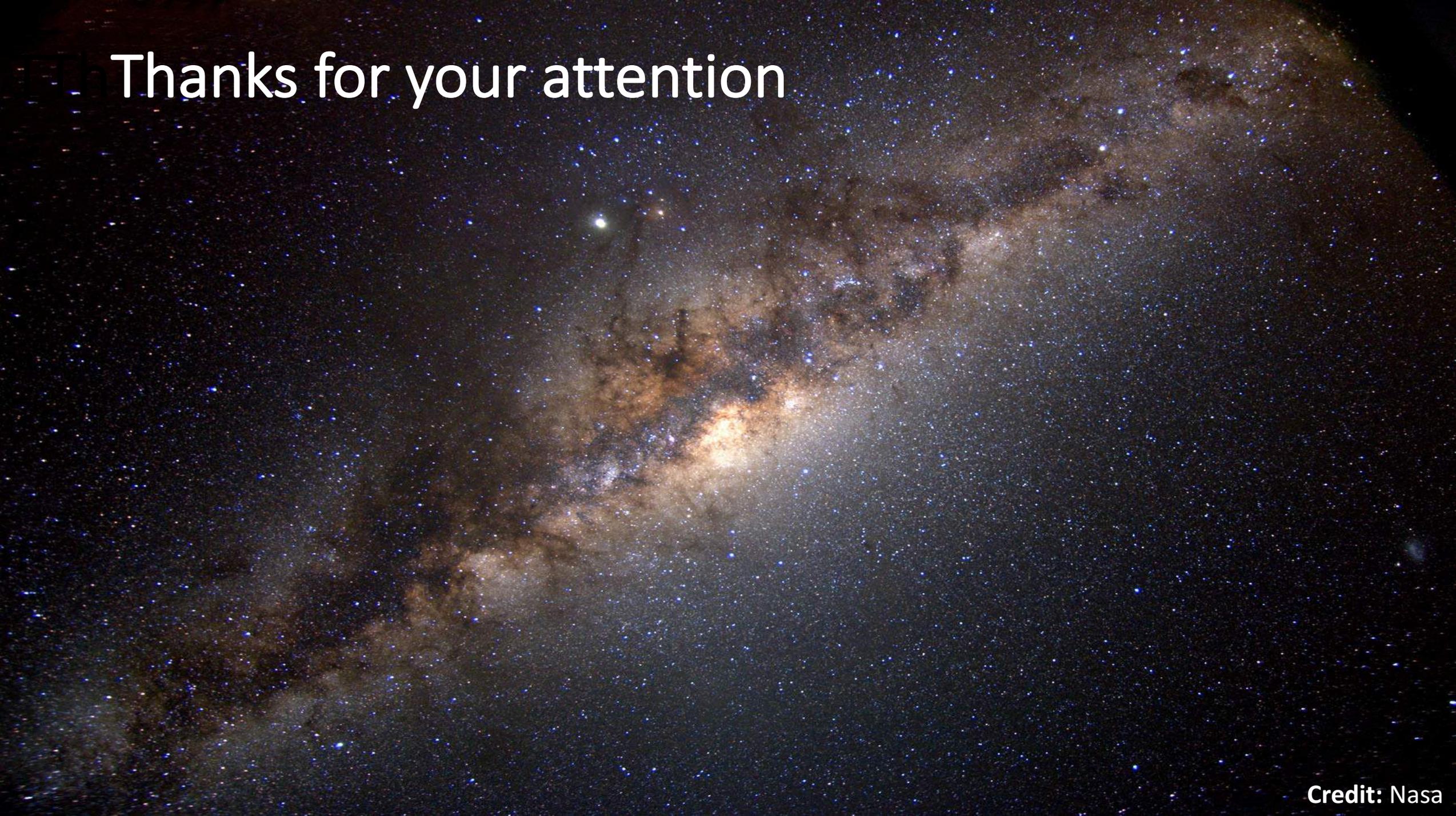


Credit: EHT Collaboration

EVOLUTION OF STARS



Thanks for your attention

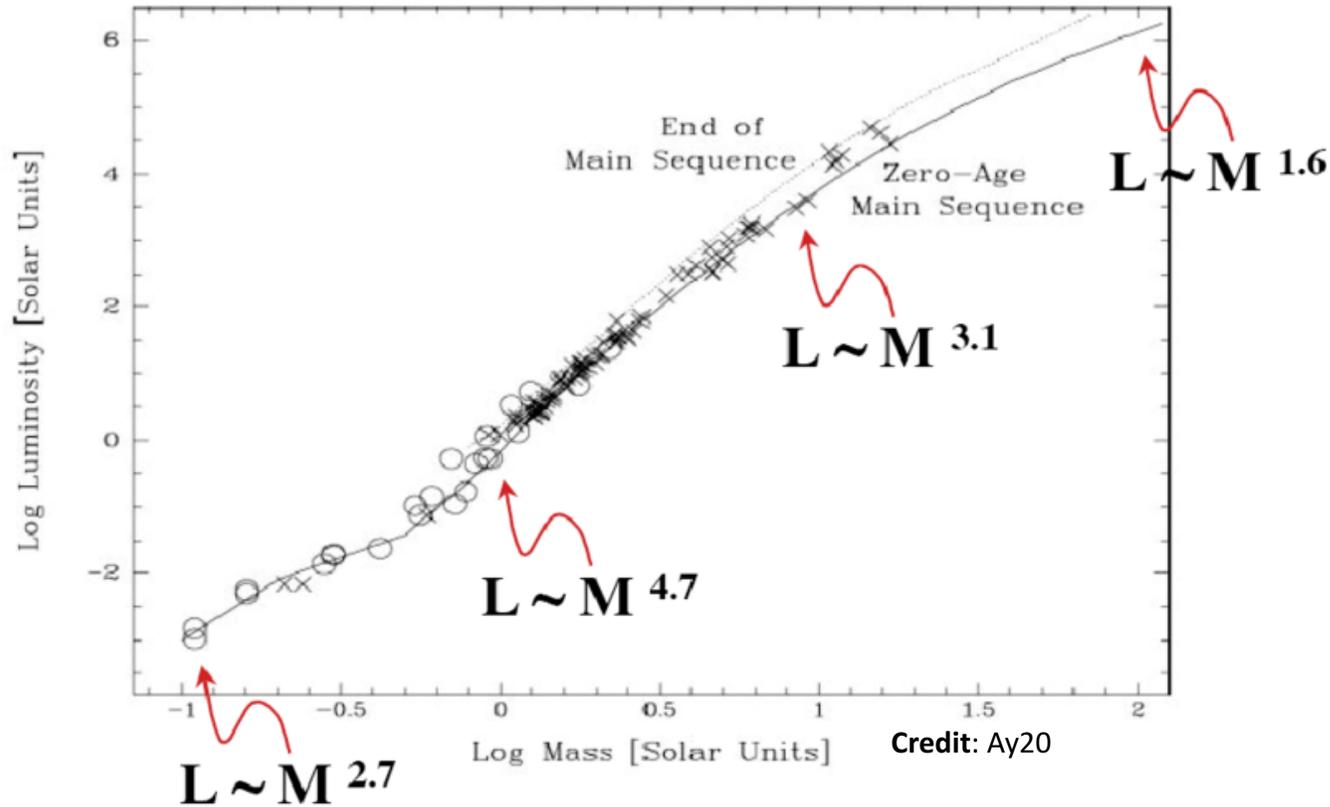


Additional Material



Main Sequence Mass-Luminosity Relation

- For MS stars a direct relationship between mass M and luminosity L is observed

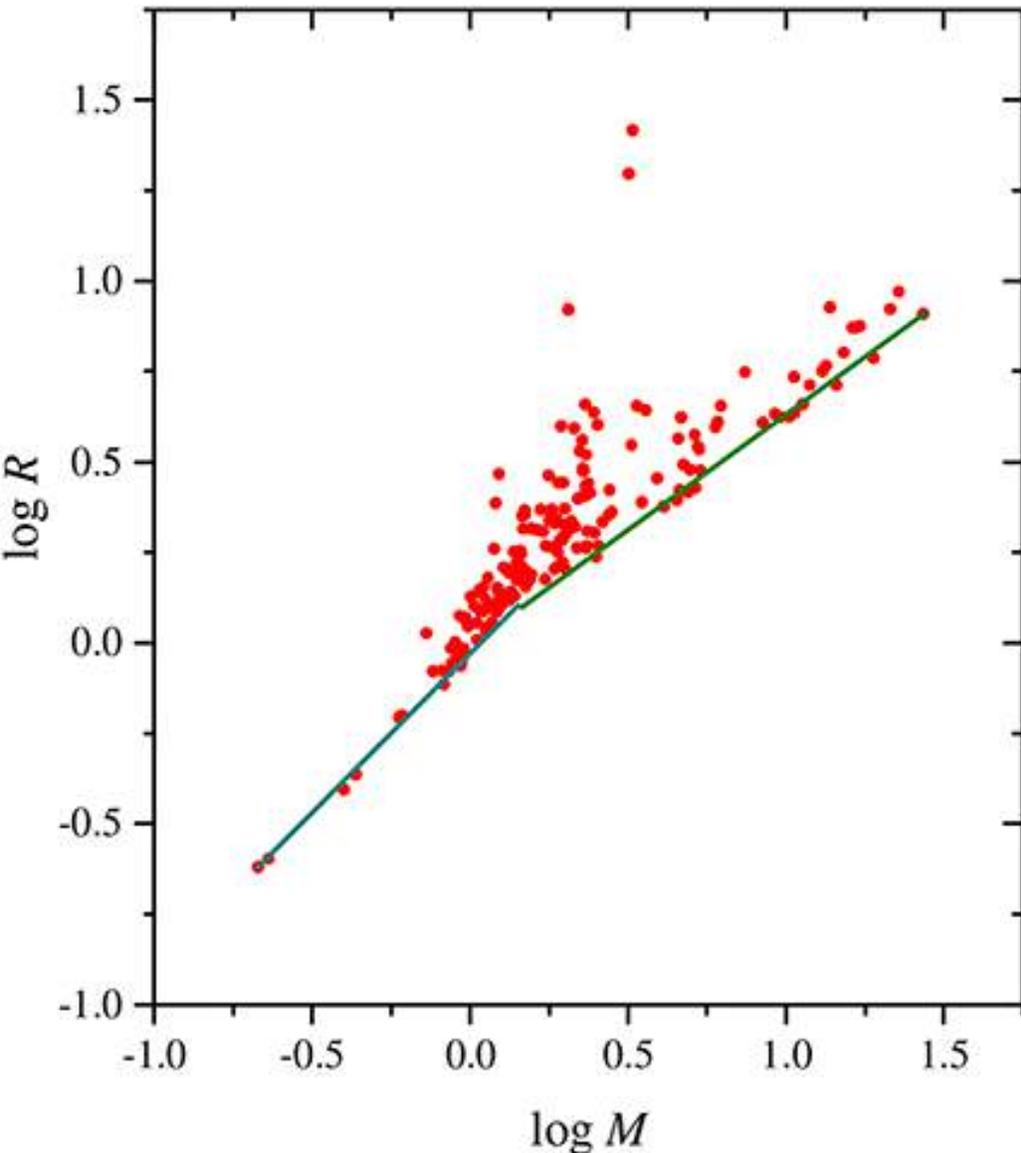


$$L \sim M^{\alpha}$$

Mass, M_{\odot}	Exponent, α
100	1,6
10	3,1
1	4,1
0,1	2,7

- it is possible to estimate the mass of a star by the observed luminosity

Main Sequence Mass-Radius Relation



- fairly tight for stars with $M < 0.9 M_{\odot}$, but large at higher mass because of stellar evolution on the Main Sequence

$$R \propto M^{\alpha}$$

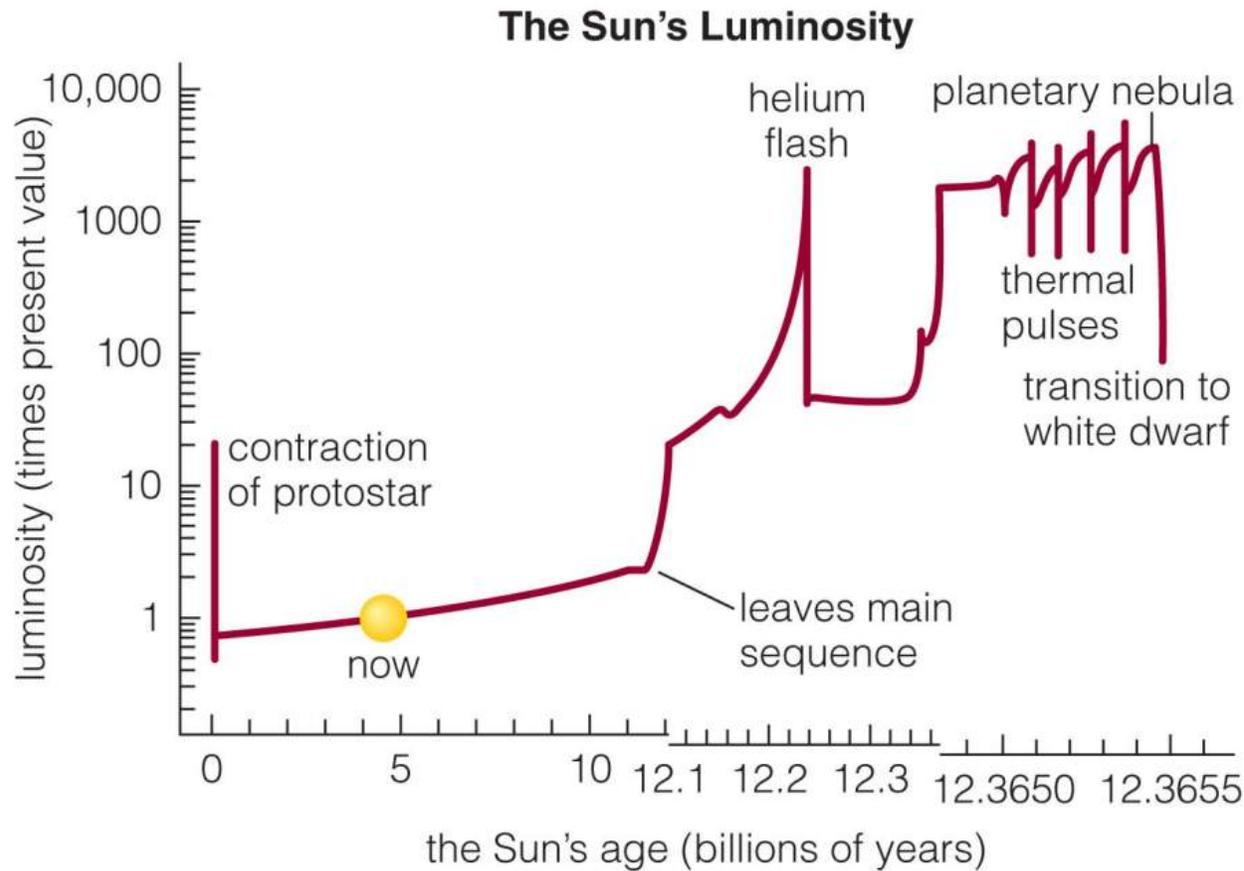
Mass, M_{\odot}	Exponent, α
0,2 – 1,4	0,9
1,4 – 30	0,6

- Mass-Radius Relation implies that

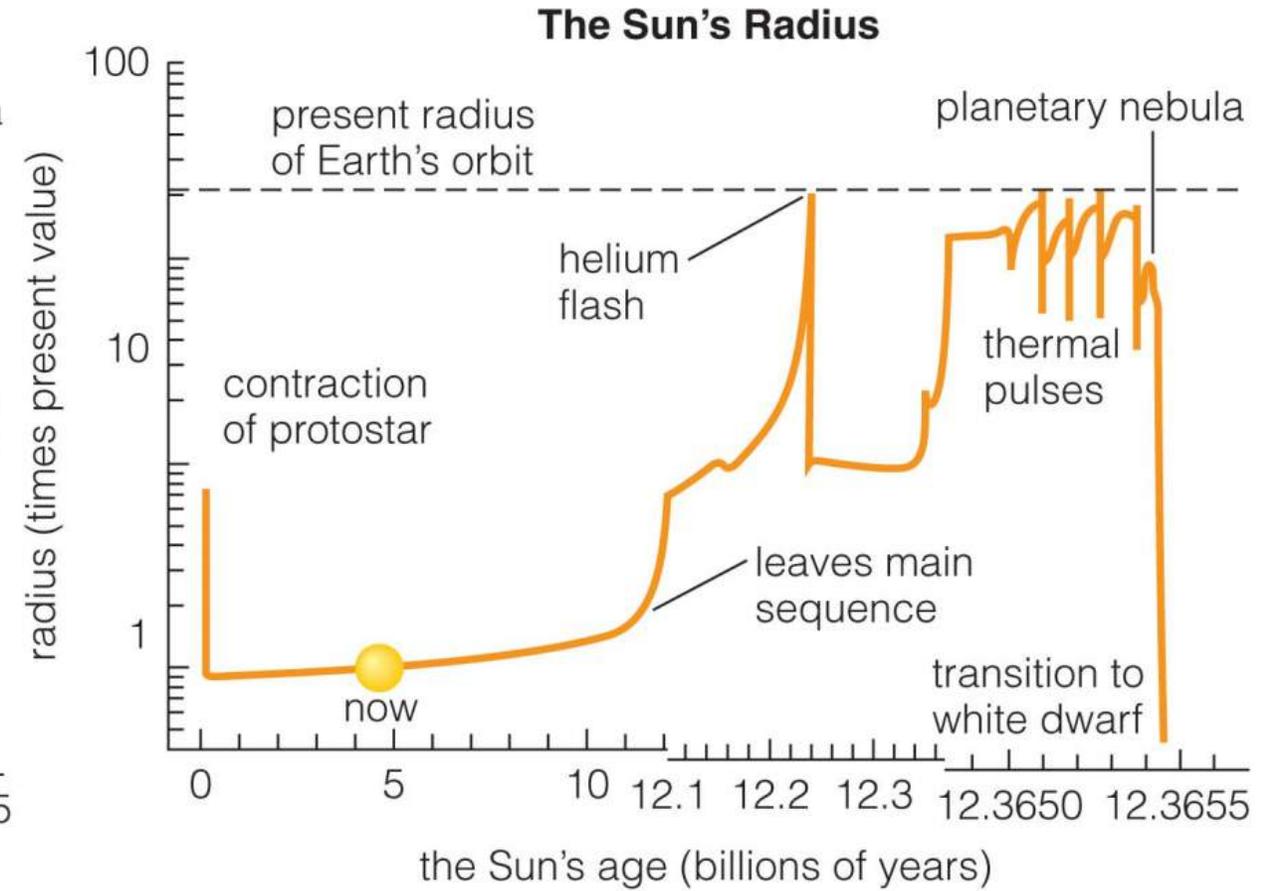
$$\bar{\rho} \propto M^{1-3\alpha}$$

→ stars are less dense with increasing mass

Comparison of Luminosity and Radius to the Lifetime of the Sun

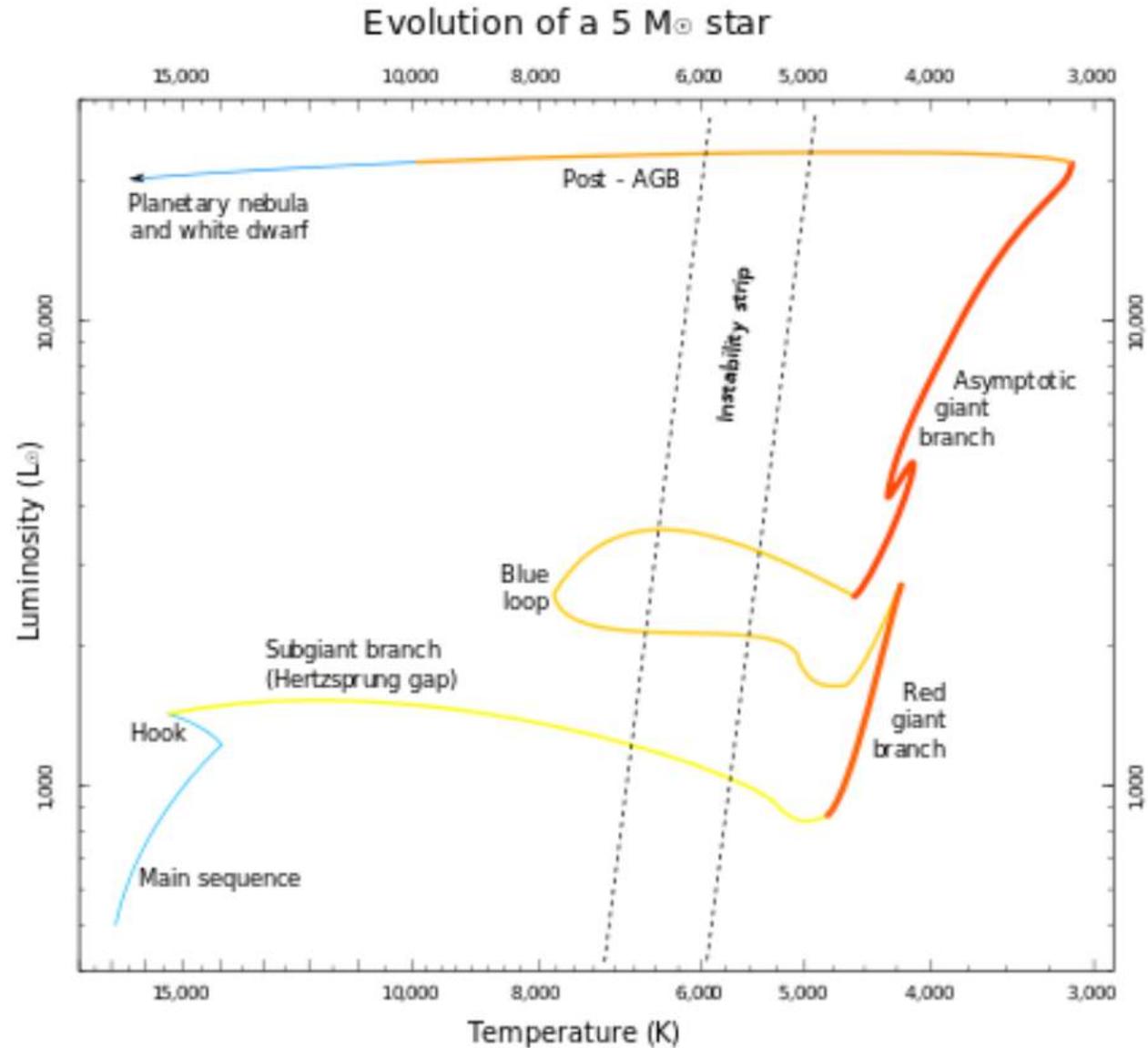
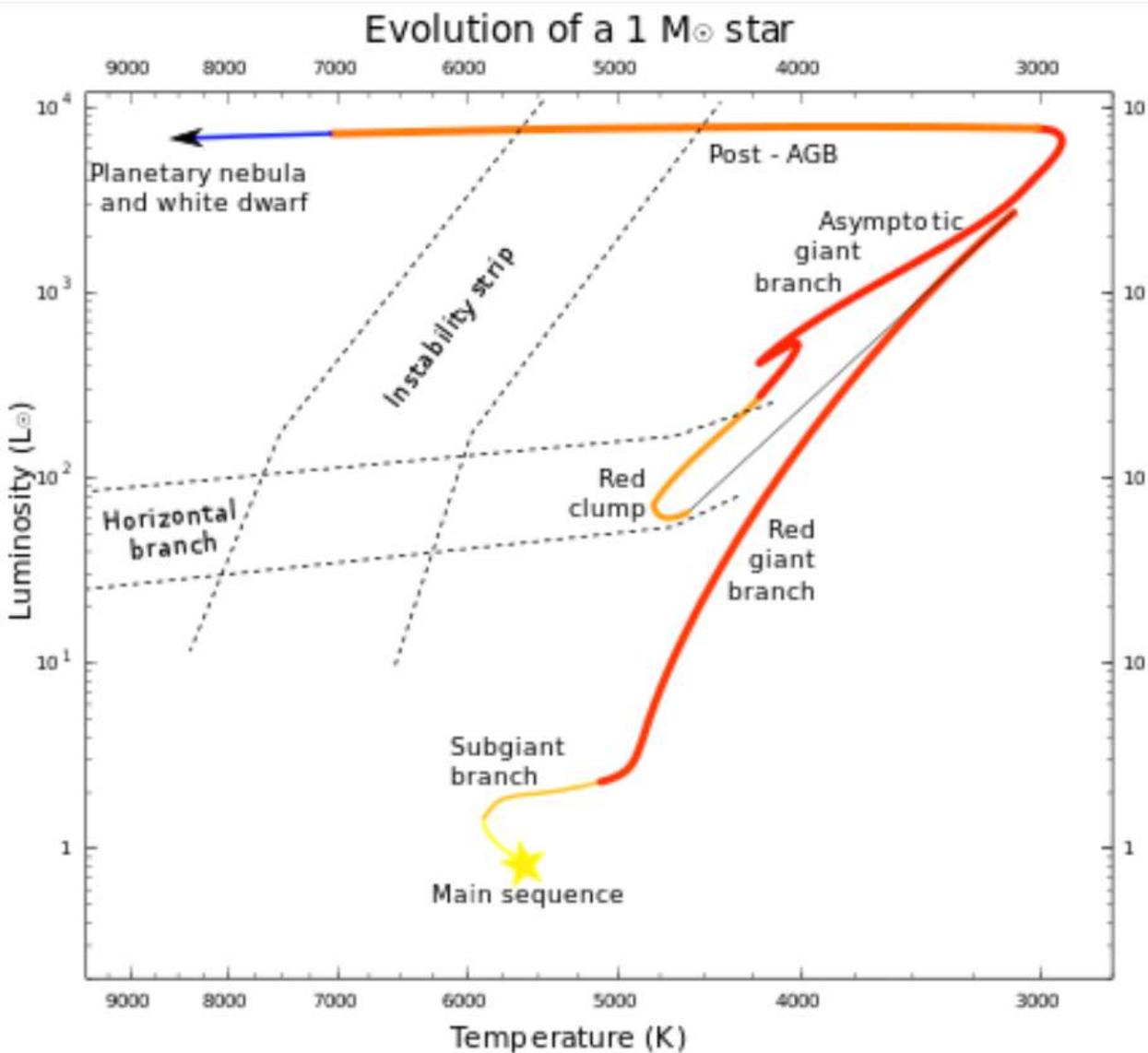


a Changes in the Sun's luminosity over time.

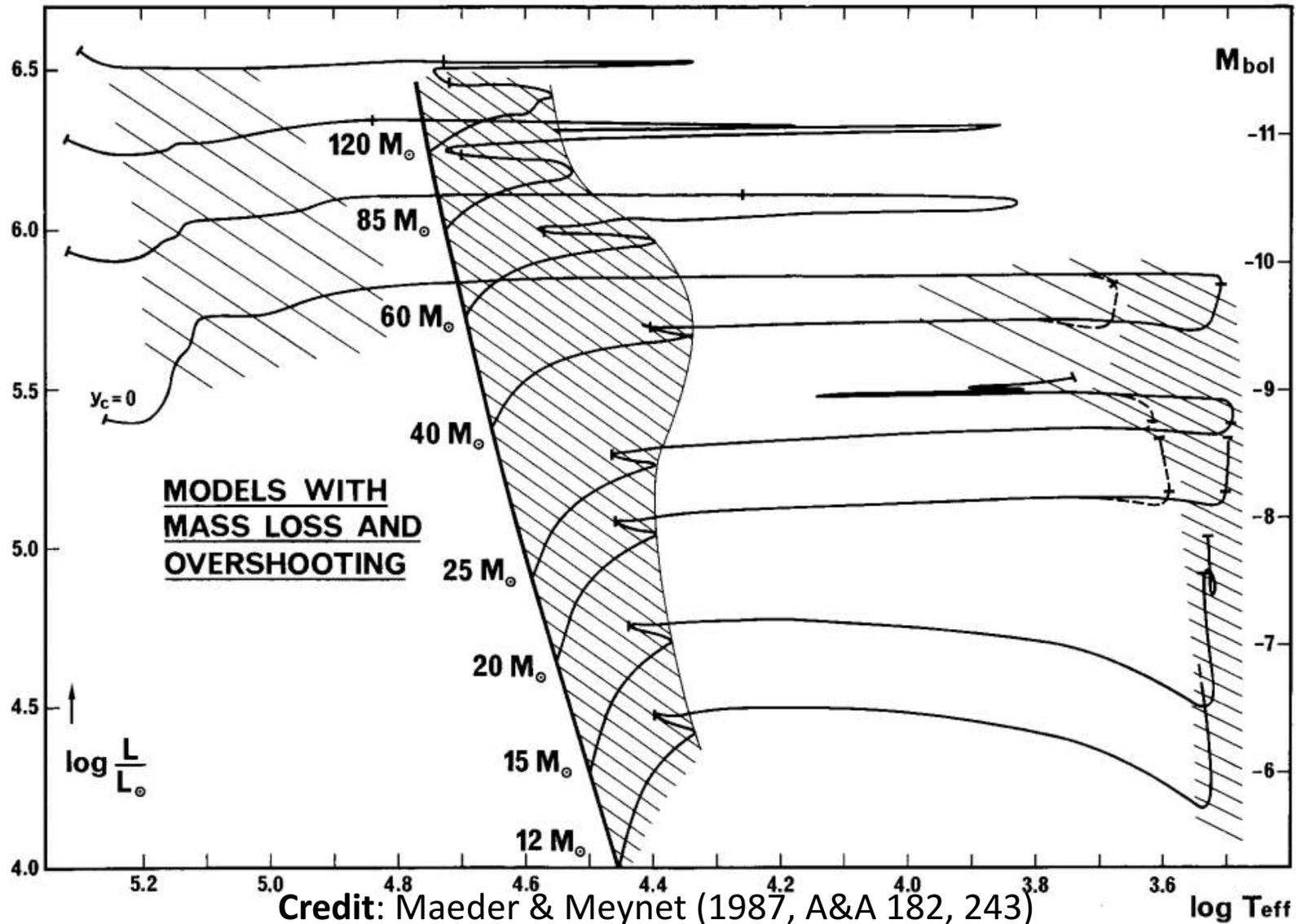


b Changes in the Sun's radius over time.

Comparison of Stellar Evolution for Low and Intermediate Mass Stars



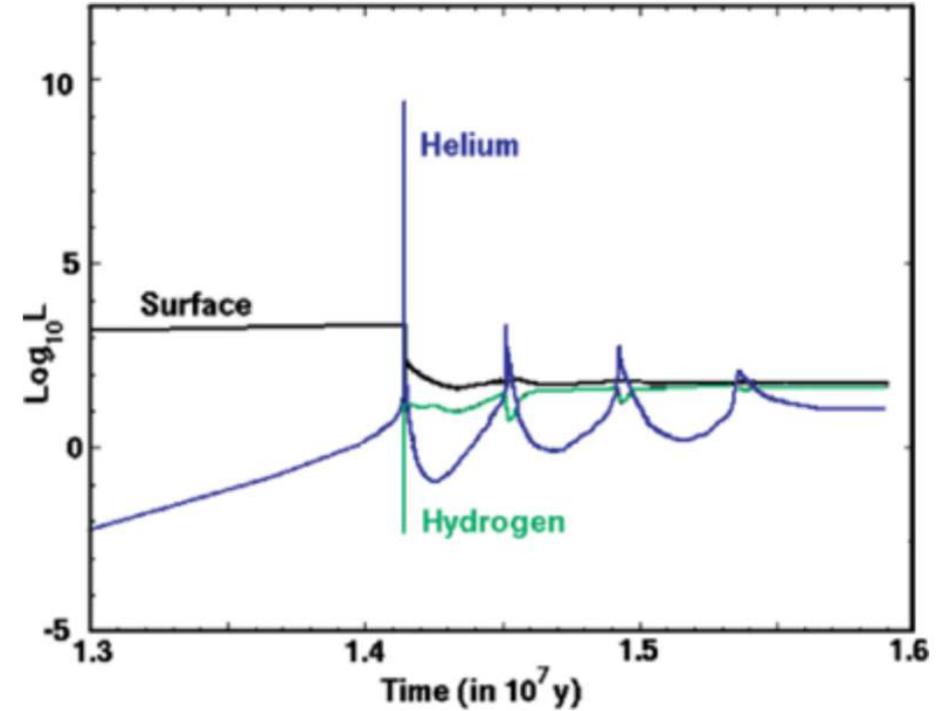
Stellar Evolution for High Mass Stars



The shaded regions correspond to long-lived evolution phases on the Main Sequence and during He core burning

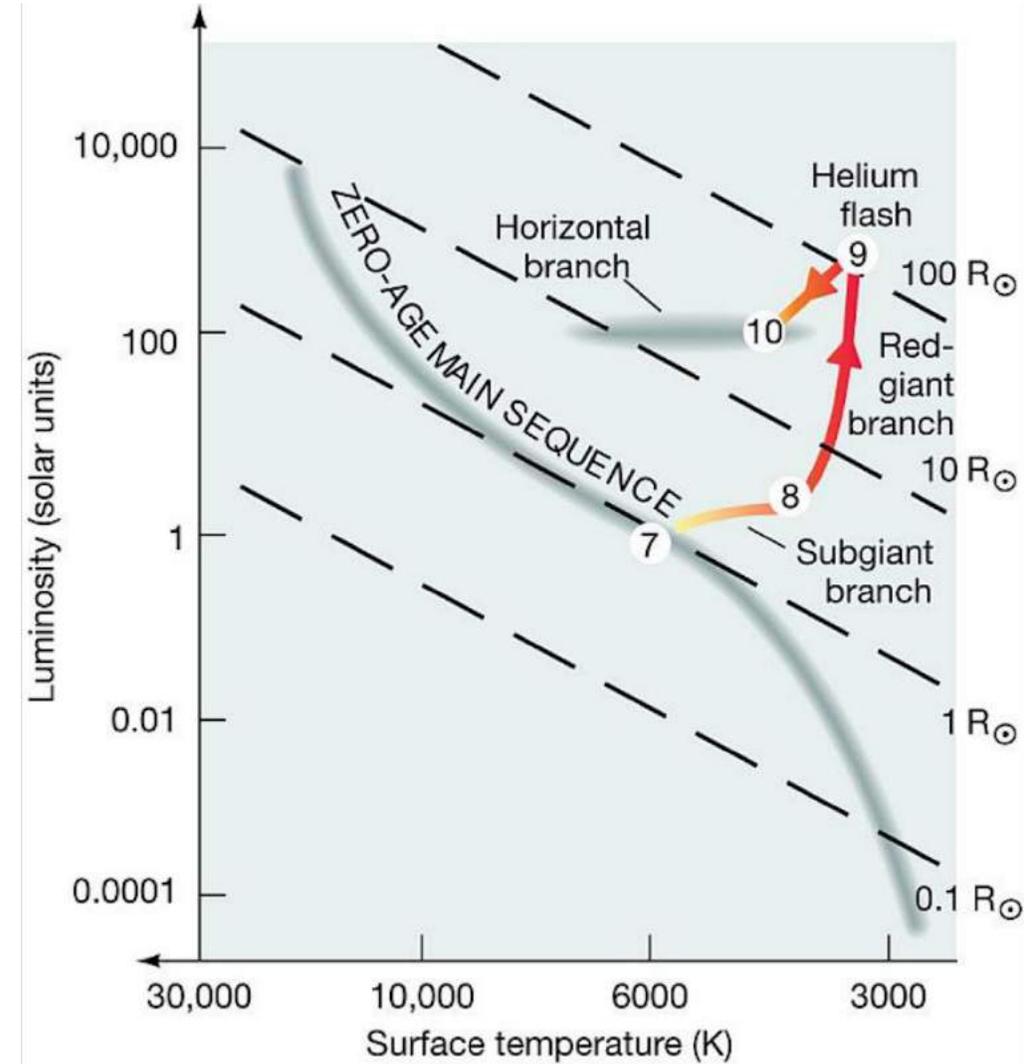
Helium Flash for $M < 2 M_{\odot}$

- stars undergo He core flash instead of regular He burning
→ cores of low mass stars are more dense than of those high mass stars
- Contracted Core is so dense that it cannot be described as an ideal gas
→ degenerate gas, in which P is only depending on ρ but not on T
- When triple- α process starts He ignites as thermonuclear runaway
→ rising T does not raise P and does not cause an expansion of the burning zone



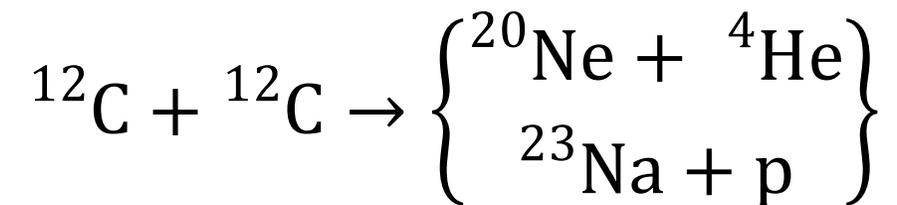
Helium Flash for $M < 2 M_{\odot}$

- mechanism of hydrostatic burning is not working anymore and He burning proceeds quickly to high T
- Very high T lift the degeneracy of the gas and its equation-of-state becomes T dependent again very suddenly
- This causes an explosive expansion of the outer core, also ejecting the outer layers of the star as a planetary nebula
→ mass loss because of He flash

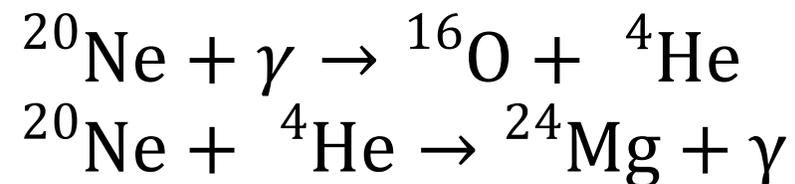


Later Burning Phases for $M > \sim 8 M_{\odot}$

- C/O core has shrunk \rightarrow core temperature increases
 \rightarrow C burning ignites at $T \sim 6 \cdot 10^8$ K

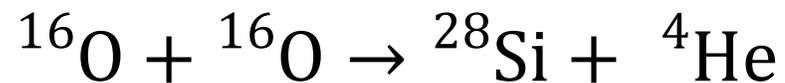


- Neon burning follows (not O!) at $T \sim 1,2 \cdot 10^9$ K
 \rightarrow photodegradation occurs for the first time and leads to the formation of heavier nuclei

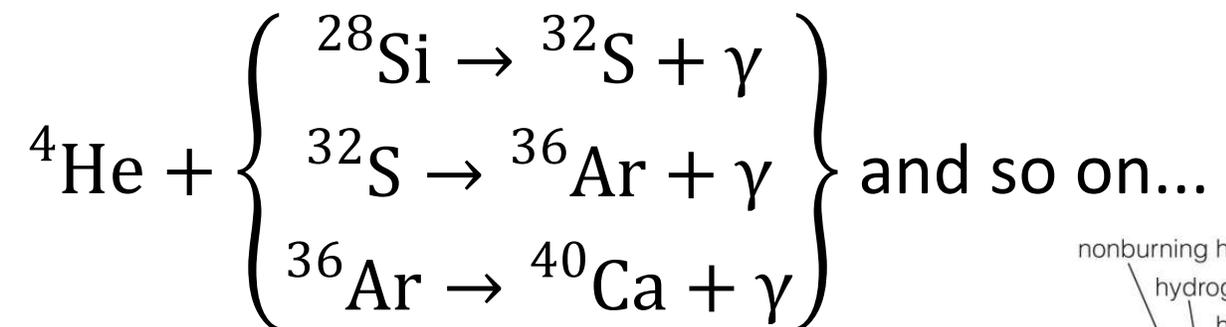


Later Burning Phases for $M > \sim 8 M_{\odot}$

- O burning occurs when temperature reaches $T \sim 1,5 \cdot 10^9$ K



- ^4He nuclei build up heavier nuclei by successive capture reactions



- Advanced burning phases lead to a „onion-skin“ layering of nuclei

- Inert Fe/Ni core has formed

